

# From AGT to knots

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# Motivation for the study of "knots"

From AGT to knots

Knots can well be the next hot topic  
in mathematical physics

Why?

Interesting

What is most exciting?

There are formulas

Many formulas



Complicated formulas  
Mysterious formulas  
Strongly interrelated formulas

Intimately related to all what we studied before:

matrix models,  
integrable systems,  
topological theories,  
CFT,  
SW theory, LMNS and Nekrasov fns, AGT relations

Related to most different parts of each story:

Unitary rather than Hermitian matrix models,  
Almost no integral formulas [but: CS models],  
only AMM/EO topological recursion [R.Dijkgraaf & H.Fuji]

Deviation from ordinary integrability:

OV pf  $\sum_R H_R \chi_R \{\bar{p}_k\}$  never ordinary  $\tau$ -functions,  
 $\sum_R \mathcal{H}_R \{p_k\} \chi_R \{\bar{p}_k\}$  are, but only sometime

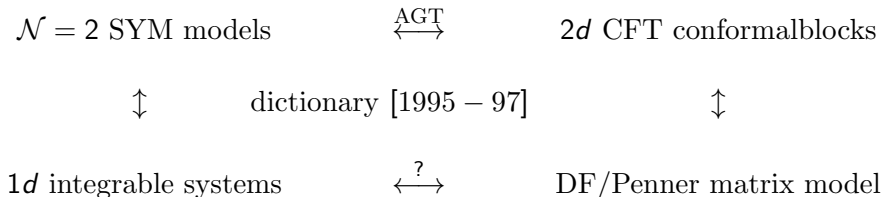
quantum  $R$ -matrices and non-trivial reps of Hecke algebras

BPS spectrum, integrality, wall crossing phenomena

modular transformations

# Universality classes are labeled by integrable systems

hep-th/9505035



quantization of integrable systems

Shroedinger-like equations (Fourier tr. of Baxter eqs.)

insertions of degenerate states

SW description through BS integrals

$$\Psi(z) = \exp \int^z \Omega, \quad \Omega = Pdz$$

$$\partial F / \partial a = \oint_B \Omega, \quad a = \oint_A \Omega$$

$$\text{NS limit } \epsilon_1 \rightarrow 0, \beta \rightarrow \infty$$

## AGT relation

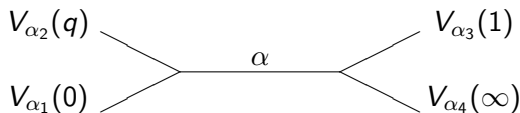
open problems  
reformulations  
interpretations  
proofs

- Dotsenko-Fateev matrix model
- Hubbard-Stratanovich duality
- Relation to integrable systems
  - Bohr-Sommerfeld integrals

One particular subject:

Beautiful  
Conceptually important  
Based on achievements in Osaka

# DF/Penner/Selberg matrix model



$$\left\langle e^{\alpha_1 \phi(0)} e^{\alpha_2 \phi(q)} e^{\alpha_3 \phi(1)} e^{\alpha_4 \phi(\infty)} \prod_{i=1}^{N_1} \int_0^q e^{b\phi(x_i)} \prod_{j=1}^{N_2} \int_0^1 e^{b\phi(y_j)} \right\rangle$$

$$\alpha_1 + \alpha_2 + bN_1 = \alpha$$

$$\alpha + \alpha_3 + \alpha_4 + bN_2 = 0$$

$$= \int dx_i \int dy_j (x_i - x_{i'})^{2\beta} (y_j - y_{j'})^{2\beta} \underline{(x_i - y_j)^{2\beta}} (x_i y_j)^{2\alpha_1 b} ((q - x_i)(q - y_j))^{2\alpha_2 b} ((1 - x_i)(1 - y_j))^{2\alpha_3 b}$$

$$= \int d\mu(x) \int d\mu(y) \left( \text{Mixing term}(x|y) \right)^2$$

# AGT as Hubbard-Stratanovich duality [1012.3137]

$$\begin{aligned}
 &\approx \int_{d\mu(x)} \int_{d\mu(y)} \exp \left( 2\beta \sum_{i,j} \log(1 - x_i y_j) \right) = \\
 &= \int_{d\mu(x)} \int_{d\mu(y)} \exp \left( \underline{2}\beta \sum_k p_k \bar{p}_k / k \right) \\
 &= \int_{d\mu(x)} \int_{d\mu(y)} \left( \sum_A \chi_A(X) \chi_A(Y) \right) \left( \sum_B \chi_B(X) \chi_B(Y) \right) \\
 &= \sum_{A,B} \left( \int_{d\mu(x)} \chi_A(X) \chi_B(X) \right) \left( \int_{d\mu(y)} \chi_A(Y) \chi_B(Y) \right)
 \end{aligned}$$

$$p_k = \text{Tr} X^k, \quad \bar{p}_k = \text{Tr} Y^k \quad [H.Itoyama \& T.Oota 1003.2929]$$

$$\exp \sum_k \frac{[\beta]_{q^k} p_k \bar{p}_k}{k} = \sum_A \frac{C_A}{C_{A'}} M_A(X) M_A(Y)$$



# AGT as Hubbard-Stratanovich duality [1012.3137]

$$\begin{array}{ccc}
 \chi_A(X) & & \chi_B(X) \\
 & \diagdown \quad \diagup & \\
 & \text{---} & \\
 & \diagup \quad \diagdown & \\
 \chi_A(Y) & & \chi_B(Y)
 \end{array}
 =
 \begin{array}{ccc}
 \chi_A(X) & & \chi_B(X) \\
 & \diagup \quad \diagdown & \\
 & | & \\
 & \diagdown \quad \diagup & \\
 \chi_A(Y) & & \chi_B(Y)
 \end{array}$$

$$\begin{aligned}
 \sum_{X,Y} \left( \sum_A \chi_A(X) \chi_A(Y) \right) \left( \sum_B \chi_B(X) \chi_B(Y) \right) &= \\
 &= \sum_{A,B} \left( \sum_X \chi_A(X) \chi_B(X) \right) \left( \sum_Y \chi_A(Y) \chi_B(Y) \right)
 \end{aligned}$$

$$\text{Conformal block} = \sum_{A,B} N_{A,B}$$

# Decomposition problem for $\beta \neq 1$

$$\int d\mu(X) \chi_A(X) \chi_B(X) \int d\mu(Y) \chi_A(Y) \chi_B(Y) \stackrel{?}{=} N_{A,B}$$

TRUE for  $\beta = 1$   
NOT so simple for  $\beta \neq 1$

$$\begin{aligned} \langle \chi_{[1]} \chi_{\bullet} \rangle \langle \chi_{[1]} \chi_{\bullet} \rangle + \langle \chi_{\bullet} \chi_{[1]} \rangle \langle \chi_{\bullet} \chi_{[1]} \rangle &= \\ &= \frac{1}{(z - \epsilon)} \frac{1}{(z + \epsilon)} + \frac{1}{(z + \epsilon)} \frac{1}{(z - \epsilon)} = \\ &= \frac{2}{z^2 - \epsilon^2} = \frac{1}{z(z - \epsilon)} + \frac{1}{z(z + \epsilon)} = N_{[1], \bullet} + N_{\bullet, [1]} \end{aligned}$$

For  $\epsilon \neq 0$  ( $\beta \neq 1$ ) particular Nekrasov functions  
have extra poles (at  $z = 0$ ), not present in Kac determinant

# Decomposition problem

Instead Nekrasov functions are nicely factorized,  
while Selberg correlators for  $\beta \neq 1$  are not:

$$\langle \chi_{[3]} \chi_{\bullet} \rangle_{BGW} \sim z^2 - (5\epsilon_1 + 8\epsilon_2)z + 6\epsilon_1^2 + 23\epsilon_1\epsilon_2 + 19\epsilon_2^2$$
$$\xrightarrow{\epsilon_2 = -\epsilon_1} z^2 + 3\epsilon_1 z + 2\epsilon_1^2 = (z + \epsilon_1)(z + 2\epsilon_1)$$

# Decomposition problem

Natural quantities, e.g. Selberg correlators  
(involved into duality relations)  
are linear combinations of  
the nicely factorized functions (Nekrasov functions),  
which possess extra singularities

AGT: generalizations

more models (quivers, 5d)

Wall crossing

hidden properties:

action of chiral algebra on the moduli spaces  
(modular transform and other ingredients of SW theory)

3d AGT

What takes the place of conformal block?

"knot polynomials"

The knot polynomials  
should be studied and understood  
at least as well as conformal blocks

Enormous amount of experimental material  
[katlas.org](http://katlas.org)

Almost no general formulas



Formulas with free parameters,  
which characterize either knots or representations,  
for:

- HOMFLY pols for generic 3, 4, (5)-strand knots in  $R = [1]$
- Colored HOMFLY for generic 3-strand braids and knots in  $R = [2], [3], [4], [5]$ 
  - Torus superpolynomials  $\mathcal{B} = [m, n]$  in  $R = [1]$   
with  $n \pmod{m} = \pm 1, \pm 2, \pm 3, \pm 4$   
(up to  $m = 9, [9, 13]$ ; the first unavailable is  $[11, 16]$ )

Mysterious hidden tropical structure is revealed in HL expansions

- Superpolynomials for the figure-eight knot  $4_1$   
in all symmetric and antisymmetric representations  $R = [p]$  and  $R = [1^p]$

# Main ideas

- Switch from knots to braids
- Introduce *extended* polynomials, depending on infinitely many time-variables
  - Study families of braids
- Derive formulas, not only pictures
  - Use matrix model techniques, (Ward identities, integrability, character expansions, integral formulas)
  - Search for universalities
  - Search for relations

$$*\mathcal{H}_R^{\mathcal{K}}(A|q)|_{A=q^N} = \left\langle \text{Tr}_R P \exp \left( \oint_{\mathcal{K}} A \right) \right\rangle_{CS}$$

Alternative definitions:

- from skein relations
- from Khovanov-Rozhansky theory –  
as Euler characteristic of the triple-graded complex

$$P(\mathbf{a}|\mathbf{q}|\mathbf{t}) = \sum_{I,J,K} N_{IJK} \mathbf{a}^I \mathbf{q}^J \mathbf{t}^K$$

$$\mathcal{H}(A|q) = P(A|q|\mathbf{t} = -1)$$

- from averages of characters  $\langle \chi_R[U] \rangle^{\mathcal{K}}$

# Hierarchy of knot polynomials for the $SL(N)$ family

For a given knot  $K$  and representation (Young diagram)  $R$

Superpolynomial  $P_R(A|q|t)$

↙  $t = q$

↘  $A = 1$

CS  $\longrightarrow$  HOMFLY  $H_R(A|q)$

Heegard – Floer  $HF_R(q|t)$

$q = 1$  ↙

$N = 2$  ↘  $N = 0$

↙  $t = q$

Special  $\sigma_R(A)$

Jones  $J_R(q)$

Alexander  $\mathcal{A}_R(q)$

# Representation dependence

$$\sigma_R(A) = \left( \sigma_{[1]}(A) \right)^{|R|}$$

$$\mathcal{A}_R(q) = \mathcal{A}_{[1]}(q^{|R|})$$

HOMFLY and superpolynomials satisfy difference equations  
as functions of the representation-variable

Garoufalidis & Le, math/0309214

Fuji, Gukov & Sulkovski, 1203.2182

IMMM, 1203.5978

## CS in temporal gauge $A_0 = 0$

$$\kappa \int \epsilon_{ijk} \text{Tr} \left( A^i \partial_j A^k + \frac{2}{3} A^i A^j A^k \right) d^3x$$
$$\xrightarrow{A_0=0} \int \text{Tr} \left( A_x \dot{A}_y dx dy \right) dt$$

Quadratic theory with the ultralocal propagator

$$\frac{2\pi i}{\kappa} \text{sign}(t) \delta(x) \delta(y)$$

Average of a Wilson line is given by  
projection onto the  $xy$  plane  
and each intersection contributes

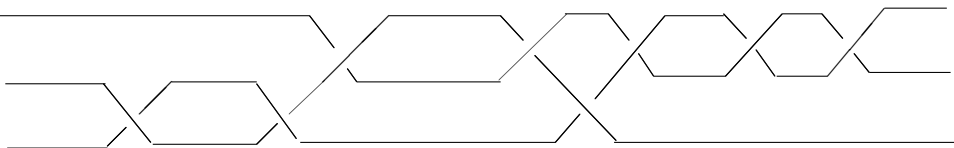
$$q^{\pm T^a \otimes T^a}$$

$$\text{with } q = e^{2\pi i / (\kappa + N)}$$

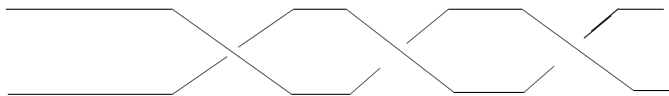
# Turaev-Reshetikhin construction

$$\mathcal{H}_R^{\mathcal{B}} = \text{Trace}_{R^{\otimes m}} \mathcal{B}$$

$$\mathcal{B} = \prod_s \mathcal{R}_{i(s), i(s)+1}^{\pm}$$



## 2-strand braids (torus knots $[2, n]$ )



$$\mathcal{H}_R^{[2,n]} = \text{Tr}_{R \otimes R} \mathcal{R}^n = \sum_Q \text{Tr}_Q \mathcal{R}^n$$

$$R \otimes R = \oplus Q$$

In each  $Q$  the  $\mathcal{R}$ -matrix acts as unity:  $\mathcal{R}_Q = \lambda_Q \cdot I_Q$

$$\mathcal{H}_R^{[2,n]} = \sum_Q \lambda_Q^n \text{Tr}_Q I_Q = \sum_Q \lambda_Q^n D_Q$$

$D_Q$  = quantum dimension of irreducible representation  $Q$



# Symmetric representations

$$[p] \otimes [p] = \bigoplus_{k=0}^p [2p - k, k]$$

$$SL(2): \text{spin } p/2 \otimes \text{spin } p/2 = \bigoplus_{k=0}^p \text{spin } (p - k)$$

$$\mathcal{R}_Q = \lambda_Q \cdot l_Q$$

$$\lambda_k = (-)^k q^{\varkappa_{[2p-k, k]}}$$

$$\varkappa_{[2p-k, k]} = 2p^2 - (2k + 1)p + k(k - 1)$$

# Universal formula

This is an example of 1-parametric general formulas,  
describing all the 2-strand braids at once

For example,  $R = [1]$ ,  $[1] \otimes [1] = [2] + [11]$ :

$$\mathcal{H}_{[1]}^{[2,n]} = q^n D_{[2]} + (-)^n q^{-n} D_{[11]}$$

Two series  $n$  odd (knots) and  $n$  even (links):

$$q^n D_{[2]} - q^{-n} D_{[11]}$$

$$q^n D_{[2]} + q^{-n} D_{[11]}$$

Also possible for other representations  
(multiparametric formula)

# W-eigenvalues

This  $\kappa_{[2p-k,k]}$  is an eigenvalue of the cut-and-join operator

$$\hat{W}[2] = \frac{1}{m} \sum_{a,b \geq 1} \left( (a+b)p_a p_b \frac{\partial}{\partial p_{a+b}} + ab p_{a+b} \frac{\partial^2}{\partial p_a \partial p_b} \right)$$

$$\hat{W}[2]S_Q\{p\} = \kappa_Q S_Q\{p\}, \quad \lambda_Q = q^{\kappa_Q}$$

$$S_1\{p\} = p_1, \quad S_2\{p\} = \frac{1}{2}(p_2 + p_1^2), \quad S_{11}\{p\} = \frac{1}{2}(-p_2 + p_1^2), \quad \dots$$

$$\kappa_Q = \sum_i q_i(q_i - 2i + 1) = \nu_Q - \nu_{Q'}$$

$$\nu_Q = \sum_i (i-1)q_i$$

$SL(N)$  characters (Shur fns)  $S_Q\{p\}$

are eigenfunctions of cut-and-join operators  $\hat{W}(\Delta)$ ,

$$\hat{W}(\Delta)S_Q = \varphi_Q(\Delta)S_Q$$

$$\hat{W}(\Delta) = : \prod_i \text{tr} \left( X \frac{\partial}{\partial X} \right)^{\delta_i} :$$

$$p_k = \text{tr} X^k = kt_k$$

These operators form a commutative algebra,  
which has a non-trivial non-commutative extension

For general theory of cut-and-join operators see [0904.4227]

# W-representation of HOMFLY for the torus knot $[2, n]$

$$\mathcal{H}_R^{[2,n]} = \sum_Q \lambda_Q^n \operatorname{Tr}_Q I_Q = \sum_Q \lambda_Q^n D_Q$$

Extended HOMFLY polynomial:

$$\mathcal{H}_R^{[2,n]} \{p_k\} = \sum_Q \lambda_Q^n S_Q \{p_k\} = q^{n\hat{W}} \sum_Q \epsilon_Q^n S_Q \{p_k\}$$

For given series ( $n$  odd or even)

$$= q^{n\hat{W}} \sum_Q \epsilon_Q S_Q \{p_k\} = q^{n\hat{W}} S_R \{p_{2k}\}$$

$$= q^{n\hat{W}} \sum_Q S_Q \{p_k\} = q^{n\hat{W}} S_R^2 \{p_k\}$$

## W-representations

Partition functions can be considered as a result of "evolution", driven by cut-and-join (W) operators from very simple "initial conditions" [0902.2627]

$$Z\{p\} = e^{g\hat{W}} \tau_0\{p\}$$

If  $W \in UGL(\infty)$ , then KP/Toda-integrability is preserved

$$\hat{W}_n = \frac{1}{2} \sum_{a,b} \left( (a+b+n)p_a p_b \frac{\partial}{\partial p_{a+b+n}} + ab p_{a+b-n} \frac{\partial^2}{\partial p_a \partial p_b} \right)$$

## W-representation. Examples

- Hermitian matrix model  $Z_N = \int dX e^{\sum_k \frac{p_k}{k} \text{Tr} X^k}$

$$Z_N = e^{\hat{W}_{-2}} e^{N p_0}$$

- Kontsevich model  $Z = \int dX e^{\text{Tr}(\frac{1}{3} X^3 - L^2 X)}$ ,  $p_k = \text{Tr} L^{-k}$

$$Z = e^{\hat{W}_{-1}^K} \cdot 1$$

$$\hat{W}_{-1}^K = \frac{2}{3} \sum (k + \frac{1}{2}) \tau_k L_{k-1}^K \quad [\text{A.Alexandrov, 1009.4887}]$$

- Hurwitz model [V.Bouchard & M.Marino, 0708.1458]

$$Z = e^{t \hat{W}_0} e^{p_1}$$

- Torus knots and links

$$Z = q^{\frac{2n}{m} \hat{W}_0} \prod_{\text{link comps}} \tilde{\chi}_R$$

Topological invariance (homotopical equivalence)  
is restored on the *topological locus*  
in the space of time variables:

$$p_k = p_k^* = \frac{A^k - A^{-k}}{q^k - q^{-k}}$$

topological invariants  $\longleftarrow$  braid invariants

$$D_Q \longleftarrow S_Q\{p_k\}: \quad S_Q\{p_k^*\}|_{A=q^N} = D_Q$$

$$D_{[1]} = [N]_q = \frac{q^N - q^{-N}}{q - q^{-1}}, \quad D_{[2]} = \frac{[N][N+1]}{[2]}, \quad D_{[11]} = \frac{[N][N-1]}{[2]}, \dots$$



# From HOMFLY to superpolynomials

$$\mathcal{H}_{[1]}^{[2,n]} = q^{n*} S_{[2]} \pm q^{-n*} S_{[11]}$$

$$\mathcal{P}_{[1]}^{[2,n]} = q^{n*} M_{[2]} \pm \left( C_{[11]} \gamma_{[11]} \right) t^{-n*} M_{[11]}$$

$$M_{[1]} \{p_{2k}\} = p_2 = M_{[2]} \{p_k\} - C_{[11]} M_{[11]} \{p_k\}$$

$$\gamma_{[11]} = \frac{1 + q^2}{1 + t^2}$$

$$\mathcal{H}_{[1,\dots,1]}^{[m,n]} \{q|p_k\} \sim L_{[m,\dots,m]} \{q^n|p_k\}$$

[1203.0667]

$$L_Q(t) = M_Q(q, t)|_{q=0}$$

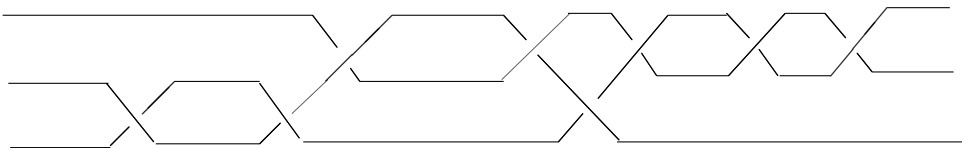
$$\mathcal{P}_{[1]}^{[m,r]} \sim \sum_{\substack{Q \vdash m \\ l(Q) \leq r}} h_Q^{(m,r)} L_Q(t)$$

$$r = 1, 2: \quad h_Q = 1$$

$$r = 3: \quad h_Q = 1 + t + (q - t)[\min(Q_1 - Q_2, Q_2 - Q_3)]_q$$

[1201.3339]

# 3-strand knots



$$\mathcal{H}_R^{(a_1, b_1 | a_2, b_2 | a_3, \dots)} = \text{Tr}_{R^{\otimes 3}} \left( \mathcal{R}_{12}^{a_1} \mathcal{R}_{23}^{b_1} \mathcal{R}_{12}^{a_2} \mathcal{R}_{23}^{b_2} \mathcal{R}_{12}^{a_3} \dots \right)$$

$$R \otimes R \otimes R = \oplus Q$$

## Reduction to the space of intertwiners [1112.2654]

$$[1] \otimes [1] \otimes [1] = ([2] + [11]) \otimes [1] = [3] + [21] + [21] + [111]$$

$$[1] \otimes [1] \otimes [1] = [1] \otimes ([2] + [11]) = [3] + [21] + [21] + [111]$$

$$\mathcal{H}_{[1]}^{(a_1, b_1 | a_2, b_2 | a_3, \dots)} = q^{a_1 + b_1 + a_2 + b_2 + \dots} S_{[3]} + (-q)^{a_1 + b_1 + a_2 + b_2 + \dots} S_{[111]} + \\ + \text{tr}_{2 \times 2} \left( \hat{\mathcal{R}}^{a_1} U \hat{\mathcal{R}}^{b_1} U^\dagger \hat{\mathcal{R}}^{a_2} U \hat{\mathcal{R}}^{b_2} U^\dagger \dots \right) S_{[21]}$$

$$\hat{\mathcal{R}} = \begin{pmatrix} q & \\ & -q^{-1} \end{pmatrix}, \quad U = \begin{pmatrix} \frac{1}{[2]} & \frac{\sqrt{[3]}}{[2]} \\ -\frac{\sqrt{[3]}}{[2]} & \frac{1}{[2]} \end{pmatrix}$$

Evaluation of HOMFLY pols  
is reduced to the study of the mixing matrices

Can they be found in general form?

# Fundamental representation $R = [1]$ , many strands $m$

$$m = 3: \quad [1]^{\otimes 3} = [3] + 2 \cdot [21] + [111]$$

$$\hat{\mathcal{R}}_2 = \begin{pmatrix} q & \\ & -\frac{1}{q} \end{pmatrix} \quad U_2 = \begin{pmatrix} c_2 & s_2 \\ -s_2 & c_2 \end{pmatrix}$$

$$m = 4: \quad [1]^{\otimes 4} = [4] + ???$$

$$\hat{\mathcal{R}}_3 = \begin{pmatrix} q & & \\ & q & \\ & & -\frac{1}{q} \end{pmatrix} \quad U_3 = \begin{pmatrix} 1 & & \\ & c_2 & s_2 \\ & -s_2 & c_2 \end{pmatrix} \quad V_3 = \begin{pmatrix} c_3 & s_3 & \\ -s_3 & c_3 & \\ & & 1 \end{pmatrix}$$

$$c_k = \frac{1}{[k]}, \quad s_k = \sqrt{1 - c_k^2} = \frac{\sqrt{[k-1][k+1]}}{[k]}$$

$m = 3$  strands, symmetric representations  $R = [\rho]$

$$\begin{pmatrix} \frac{\sqrt{\lambda\mu}}{\lambda+\mu} & \frac{\sqrt{\lambda^2+\lambda\mu+\mu^2}}{\lambda+\mu} \\ -\frac{\sqrt{\lambda^2+\lambda\mu+\mu^2}}{\lambda+\mu} & \frac{\sqrt{\lambda\mu}}{\lambda+\mu} \end{pmatrix}$$

$$\begin{pmatrix} -\frac{\lambda_1(\lambda_2-\mu)}{(\lambda_1-\lambda_2)(\lambda_1+\mu)} & \frac{\sqrt{(\lambda_1\lambda_2+\mu^2)(\lambda_1^2-\mu\lambda_2)}}{\sqrt{(\lambda_1-\lambda_2)(\lambda_2+\mu)}(\lambda_1+\mu)} & -\frac{1}{\lambda_1-\lambda_2} \sqrt{\frac{(\mu\lambda_1-\lambda_2^2)(\lambda_1^2-\mu\lambda_2)}{(\lambda_1+\mu)(\lambda_2+\mu)}} \\ -\frac{\sqrt{(\lambda_1\lambda_2+\mu^2)(\lambda_1^2-\mu\lambda_2)}}{\sqrt{(\lambda_1-\lambda_2)(\lambda_2+\mu)}(\lambda_1+\mu)} & -\frac{(\lambda_1+\lambda_2)\mu}{(\lambda_1+\mu)(\lambda_2+\mu)} & -\frac{\sqrt{(\lambda_1\lambda_2+\mu^2)(\mu\lambda_1-\lambda_2^2)}}{\sqrt{(\lambda_1-\lambda_2)(\lambda_1+\mu)}(\lambda_2+\mu)} \\ -\frac{1}{\lambda_1-\lambda_2} \sqrt{\frac{(\mu\lambda_1-\lambda_2^2)(\lambda_1^2-\mu\lambda_2)}{(\lambda_1+\mu)(\lambda_2+\mu)}} & \frac{\sqrt{(\lambda_1\lambda_2+\mu^2)(\mu\lambda_1-\lambda_2^2)}}{\sqrt{(\lambda_1-\lambda_2)(\lambda_1+\mu)}(\lambda_2+\mu)} & \frac{\lambda_2(\lambda_1-\mu)}{(\lambda_1-\lambda_2)(\lambda_2+\mu)} \end{pmatrix}$$

...

## Trefoil and the figure eight knot

$$\text{trefoil: } 3_1 = [2, 3] = [3, 2] = (1, 1|1, 1)$$

$$\mathcal{H}_{[1]}^{3_1} = q^3 S_{[2]} - q^{-3} S_{[11]}$$

$$\begin{aligned} \mathcal{H}_{[1]}^{3_1} &= q^4 S_{[3]} + q^{-4} S_{[111]} + \text{tr}_{2 \times 2} \left( \hat{R} U \hat{R}^{\pm 1} U^\dagger \hat{R} U \hat{R}^{\pm 1} U^\dagger \dots \right) S_{[21]} = \\ &= \underbrace{\frac{A - A^{-1}}{q - q^{-1}}}_{*S_{[1]}(A|q)} \left( q^4 A - (q^2 + q^{-2}) + q^{-4} A^{-4} \right) \end{aligned}$$

$$\text{figure eight knot: } 4_1 = (1, -1|1, -1)$$

$$\mathcal{H}_{[1]}^{4_1} = S_{[3]} + S_{[111]} + \text{tr}_{2 \times 2} \left( \hat{R} U \hat{R}^{-1} U^\dagger \hat{R} U \hat{R}^{-1} U^\dagger \dots \right) S_{[21]}$$



# Application to the figure eight knot [1203.5978]

$$\begin{aligned}\frac{{}^*\mathcal{H}_{[1]}^{4_1}(A|q)}{{}^*S_{[1]}(A|q)} &= A^2 - (q^2 - 1 + q^{-2}) + A^{-2} = \\ &= 1 + (Aq - (Aq)^{-1})(Aq^{-1} - A^{-1}q^{-1}) = 1 + \{Aq\}\{Aq^{-1}\} \\ &\quad \{x\} = x - x^{-1} \\ [1] \times [1] \times [1] &= [3] + 2 \cdot [21] + [111]\end{aligned}$$

$$[2] \times [2] \times [2] = [6] + 2 \cdot [51] + 3 \cdot [42] + [411] + [33] + 2 \cdot [321] + [222]$$

$$\frac{{}^*\mathcal{H}_{[2]}^{4_1}(A|q)}{{}^*S_{[2]}(A|q)} = 1 + [2]_q \{Aq\}\{Aq^{-1}\} + \{Aq^3\}\{Aq^2\}\{A\}\{Aq^{-1}\}$$

[P.Ramadevi and T.Sarkar, hep-th/0009188]

## Classical ( $q = 1$ ) case

$$\sigma_R(A) = \lim_{q=1} \frac{{}^* \mathcal{H}_R^{4_1}(A|q)}{{}^* \mathcal{S}_R(A|q)} = \left(\sigma_{[1]}(A)\right)^{|R|}$$

$$\sigma_{[1]}(A)^{4_1} = 1 + \{Aq\}\{Aq^{-1}\}|_{q=1} = 1 + \{A\}^2 = 1 + (A - A^{-1})^2$$

$$\sigma_R^{4_1}(A) = \left(1 + \{A\}\right)^{|R|} = \sum_{k=0}^{|R|} C_k^{|R|} \{A\}^{2k}$$

$$\sigma_R^{4_1}(A) = \left(1 + \{A\}\right)^{|R|} = \sum_{k=0}^{|R|} \frac{|R|!}{k!(|R| - k)!} \{A\}^{2k}$$

$$\frac{{}^*\mathcal{H}_{[1]}^{4_1}(A|q)}{{}^*\mathcal{S}_{[1]}(A|q)} = 1 + \{Aq\}\{Aq^{-1}\}$$

$$\frac{{}^*\mathcal{H}_{[2]}^{4_1}(A|q)}{{}^*\mathcal{S}_{[2]}(A|q)} = 1 + [2]_q \{Aq\}\{Aq^{-1}\} + \{Aq^3\}\{Aq^2\}\{A\}\{Aq^{-1}\}$$

...

$$\frac{{}^*\mathcal{H}_{[p]}^{4_1}(A|q)}{{}^*\mathcal{S}_{[2]}(A|q)} = \sum_{k=0}^p \frac{[p]!}{[k]![p-k]!} \prod_{i=0}^{p-1} \{Aq^{p+i}\}\{Aq^{i-1}\}$$

We have seven pieces of evidence:

Our answer

- Reproduces particular examples at  $R = [2], [3], [4], [5]$
- For  $q \rightarrow 1$  reproduces the conjectured special polynomials
- Consistent with the interesting formula, describing the value of  ${}^* \mathcal{H}_{[p]}^{4_1}(A|q)$  at the one-dimensional locus  $q = e^{\frac{i\pi}{N+p-1}}$ ,  $A = q^N = -e^{\frac{i\pi(1-p)}{N+p-1}}$
- For  $A = q^2$  reproduces the known answers for the Jones polynomials
  - For  $A = 1$  reproduces the Alexander polynomial
- Related antisymmetric HOMFLY polynomial  ${}^* \mathcal{H}_{[1^p]}^{4_1}(A|q)$  vanishes for  $A = q^N$  with  $N < p$ , i.e. whenever  $p$  exceeds the rank of the group by two, and turns its ratio to the unknot turns into unity for  $N = p$ .
  - Consistent with the Ooguri-Vafa conjecture

# Superpolynomial for $4_1$

$$\sum_{k=0}^p \frac{[|R|]!}{[k]![|R| - k]!} \prod_{i=1}^k Z_i(A) = \sum_{k=0}^p \prod_{i_1 \leq \dots \leq i_k} Z_{i_1}(A) Z_{i_2}(Aq) Z_{i_3}(Aq^2) \dots Z_{i_k}(Aq^{k-1})$$

$$Z_i(A) = \{Aq^{2(p-i)+1}\} \{Aq^{-1}\} \longrightarrow \mathfrak{Z}_i(A) = \{Aq^{2(p-i)+1}\} \{At^{-1}\}$$

$$\frac{{}^* \mathcal{P}_{[p]}^{4_1}(A|q, t)}{{}^* M_{[p]}(A|q, t)} = \sum_{k=0}^p \prod_{1 \leq i_1 \leq \dots \leq i_k \leq p} \mathfrak{Z}_{i_1}(A) \mathfrak{Z}_{i_2}(Aq) \mathfrak{Z}_{i_3}(Aq^2) \dots \mathfrak{Z}_{i_k}(Aq^{k-1})$$

$$t = \mathbf{q}, \quad q = -\mathbf{qt}, \quad A = \mathbf{a}\sqrt{-\mathbf{t}}$$

$$\mathfrak{Z}_i(Aq^s) = \frac{(1 + \mathbf{a}^2 \mathbf{t} (\mathbf{qt})^{4(p-i)+2+2s}) (\mathbf{q}^2 + \mathbf{a}^2 \mathbf{t} (\mathbf{qt})^{2s})}{\mathbf{a}^2 \cdot (\mathbf{qt})^{2(p-i+s+1)}}$$

# Difference equation

$$P_{[\rho+1]}(A) - P_{[\rho]}(A) = \{Aq^{2\rho+1}\}\{At^{-1}\}P_{[\rho]}(qA)$$

$$P_{[\rho]} = \frac{{}^*\mathcal{P}_{[\rho]}^{4_1}(A|q, t)}{{}^*M_{[\rho]}(A|q, t)}$$

# Generalizations

- from  $R = [p]$  to arbitrary  $R$  (arbitrary Young diagram

$$R = \{p_1 \geq p_2 \geq \dots \geq 0\}$$

- from  $\mathcal{K} = (1, -1|1, -1)$  to entire series of 3-strand knots

$\mathcal{K} = (1, -1)^n = (1, -1|1, -1|\dots|1, -1)$ , a simple generalization of the torus knot family  $\mathcal{K} = [3, n] = (1, 1)^n$  (these are knots for  $n$ , indivisible by  $m = 3$  and 3-component links otherwise)

MANY THANKS FOR YOUR ATTENTION!



THANKS TO THE ORGANIZERS!!!